



Interstellar Neutral Helium in the Heliosphere from *IBEX* Observations. V. Observations in *IBEX*-Lo ESA Steps 1, 2, and 3

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Abstract

Direct-sampling observations of interstellar neutral (ISN) He by the *Interstellar Boundary Explorer* (*IBEX*) provide valuable insight into the physical state of and processes operating in the interstellar medium ahead of the heliosphere. The ISN He atom signals are observed at the four lowest ESA steps of the *IBEX*-Lo sensor. The observed signal is a mixture of the primary and secondary components of ISN He and H. Previously, only data from one of the ESA steps have been used. Here, we extend the analysis to data collected in the three lowest ESA steps with the strongest ISN He signal, for the observation seasons 2009–2015. The instrument sensitivity is modeled as a linear function of the atom impact speed onto the sensor's conversion surface separately for each ESA step of the instrument. We find that the sensitivity increases from lower to higher ESA steps, but within each of the ESA steps it is a decreasing function of the atom impact speed. This result may be influenced by the hydrogen contribution, which was not included in the adopted model, but seems to exist in the signal. We conclude that the currently accepted temperature of ISN He and velocity of the Sun through the interstellar medium do not need a revision, and we sketch a plan of further data analysis aiming at investigating ISN H and a better understanding of the population of ISN He originating in the outer heliosheath.

Key words: instrumentation: detectors – ISM: atoms – ISM: kinematics and dynamics – local interstellar matter – methods: data analysis – Sun: heliosphere

1. Introduction

Direct sampling of interstellar neutral (ISN) gas in the heliosphere is a powerful tool for investigating the physical state of the interstellar matter ahead of the heliosphere (Möbius et al. 2009). The ISN gas is a mixture of various elements, mostly hydrogen and helium (Frisch et al. 2011). Strong ionization processes deplete the population of ISN H close to the Sun and thus ISN He is the most abundant species at 1 au (Ruciński et al. 2003). Due to the long mean free path of the neutral He atoms in the interstellar medium most of the observed ISN atoms are created in the unperturbed medium well ahead of the heliosphere (Bzowski et al. 2017). These atoms are collectively called the primary ISN population. Moreover, an additional population of atoms is created in the interstellar medium modified due to interaction with the heliosphere—the secondary ISN population (Baranov & Malama 1995).

Results of in-depth analyses of direct-sampling observations from *GAS/Ulysses* (Witte 2004; Bzowski et al. 2014; Wood et al. 2015) and from *IBEX*-Lo (Bzowski et al. 2015; McComas et al. 2015a, 2015b; Möbius et al. 2015a; Schwadron et al. 2015) provided the direction, speed, and temperature of the primary ISN He inflowing to the heliosphere. In addition to the primary ISN He, *IBEX*-Lo discovered the Warm Breeze (Bzowski et al. 2012; Kubiak et al. 2014). This most likely is the secondary population of ISN He, created via charge exchange between interstellar He⁺ ions and ISN He atoms in the outer heliosheath (Kubiak et al. 2016; Bzowski et al. 2017) and thus bears information on the plasma flow in the region

beyond the heliopause. The signal measured by *IBEX*-Lo cannot be easily separated into the primary and secondary components. Moreover, due to the measurement technique used by *IBEX*-Lo (Möbius et al. 2009; Park et al. 2015), it is challenging to uniquely identify the counts registered by the *IBEX*-Lo instrument due to neutral He and H atoms. Consequently, the signal observed by *IBEX*-Lo is a sum of the primary and secondary populations of ISN He and ISN H.

ISN H has been identified in *IBEX*-Lo observations (Saul et al. 2012, 2013; Schwadron et al. 2013), but the understanding of the signal in different ESA steps of *IBEX*-Lo is still not satisfactory (Schwadron et al. 2013; Katushkina et al. 2015). One of the reasons was an inaccurate knowledge of the radiation pressure acting on H atoms in the heliosphere (Kowalska-Leszczynska et al. 2018). The ability to resolve the ISN H and ISN He components in the observed signal is important because ISN H plays a vital role in general heliospheric studies and because neglecting the hydrogen contribution in the He signal may bias the inferred parameters of inflow of interstellar matter on the heliosphere, as illustrated by the differences in the estimates of these parameters by Bzowski et al. (2012) and Möbius et al. (2012) on one hand and by Bzowski et al. (2015), Leonard et al. (2015), Möbius et al. (2015a), and Schwadron et al. (2015) on the other. A discussion of this topic was provided by Swaczyna et al. (2015) and Bzowski et al. (2015).

IBEX-Lo is a time-of-flight mass spectrometer (Fuselier et al. 2009) that measures neutral atoms by registering their negative ions, created upon impact on a specially prepared conversion

surface. The conversion surface is permanently covered with a thin layer of water, constantly replenished due to outgassing of material from the sensor and the spacecraft. Thus, in addition to negative ions from the direct ionization mechanism, negatively charged products of sputtering of material from the conversion surface are registered (Wurz et al. 2006). In the case of H atoms, the first mechanism dominates. However, He atoms rarely form negative ions (Wurz et al. 2008) and therefore they are detected owing to the second mechanism. The products sputtered by He atoms from the water layer make the *IBEX*-Lo instrument sensitive to He atoms with energies typical for ISN atoms at 1 au from the Sun. The ions sputtered by He atoms from this water layer are primarily H^- and O^- ions. Thus, the signal interpreted as coming from He atoms can originate both from H^- ions sputtered off the conversion surface by the incoming He atoms and from real H atoms reflected and ionized on the conversion surface. Disentangling the two contributions requires a detailed analysis (Park et al. 2016). Calibration of the sensitivity of *IBEX*-Lo to H atoms with various energies was carried out in the laboratory, but this was possible only within certain limitations due to challenges in obtaining a monoenergetic beam of neutral atoms with relevant energies (Wieser & Wurz 2005) and because the water layer on the conversion surface may be different between the laboratory and space conditions.

IBEX-Lo does not directly measure the energy of the impacting atoms. It detects ions created at the conversion surface with energies within the energy range set by the electrostatic analyzer (ESA) and rejects those with the energies outside this preselected range. Ions with different energies are observed when the instrument is switched to a different energy setting. The atoms directly converted to negative ions preserve most of their energies, but the ions sputtered from the conversion surface have a broad energy distribution, limited by the energy of the impacting atoms. As discussed later in the paper, ISN He is visible in the four lowest energy settings of the ESA (hereafter ESA steps), but so far only measurements from ESA step 2 have been used in the analyses of ISN He. In this paper, we extend the analysis to ESA steps 1 and 3. ESA step 4 is left out because of significantly lower count rates (by a factor of ~ 5). A better insight into the physics of interaction of the heliosphere with interstellar matter could potentially be obtained if data from the three ESA steps with the highest count rates are used. The factor that has been preventing the use of data from all of them is the lack of relative calibration of the ESA steps for He atoms.

In the preliminary analyses (Möbius et al. 2009, 2012; Bzowski et al. 2012) it was assumed that the sensitivity function to various energies in ESA step 2 is flat, i.e., that the sensitivity is independent of energy. However, Kubiak et al. (2014), Sokół et al. (2015a), and Galli et al. (2015) showed evidence that, most likely, there is a threshold for the sensitivity to He atoms at an energy between 19 and 38 eV. This is low compared to the ~ 130 eV incident energy of the ISN He atoms. However, this threshold was shown to be particularly important for the analysis of the Warm Breeze (Kubiak et al. 2014, Figure 8) and for the attempt to find ISN He in observations carried out during the fall ISN observation seasons (Galli et al. 2015).

In this paper we make the first attempt to simultaneously determine the inflow velocity vector, the temperature of ISN He, and the relative energy sensitivity characteristics of the *IBEX*-Lo detector to He atoms using data from the three lowest

ESA steps. We use measurements from the ISN observation campaigns 2009 through 2015, i.e., one season more than in the previous analyses (Bzowski et al. 2015; Möbius et al. 2015a). First, we discuss the data, the adopted models, and the analysis method (Section 2). Then, we present results of the analysis, including the relative energy sensitivity characteristics of the three ESA steps of *IBEX*-Lo (Section 3). Finally, we discuss some implications of the findings (Section 4).

2. Methods

2.1. *IBEX*-Lo Observations

IBEX-Lo observations are carried out in eight partially overlapping, logarithmically spaced, sequentially switched energy channels (the ESA steps). The full width at half maximum of the ESA steps is $\Delta E \simeq 0.7E$, where E is the central energy of a given ESA step. For our purposes, the most relevant are the three lowest channels, i.e., ESA steps 1–3, with central energies 15, 29, and 55 eV, respectively (Fuselier et al. 2009, 2012). The absolute sensitivity of the instrument depends on the post-acceleration (PAC) voltage in the ESA. This value is generally kept constant, but once during the mission, after the fourth ISN campaign, it was reduced. As discussed by Bzowski et al. (2015), this resulted in an approximately twofold reduction in the absolute sensitivity to He atoms. The central energies and widths of the ESA steps correspond to energies of incident H atoms that are converted to negative ions. The energy loss is small in this situation. Helium atoms are observed due to the sputtering of material from the conversion surface. This process results in a significantly larger energy difference between the incident He atom and the sputtered H^- ion. The sputtered ions have a wide energy spectrum and are observed in the ESA steps with central energies much below the energies of the incident ISN He atoms.

Details of the operation of the *IBEX*-Lo instrument during the yearly campaigns of observations of ISN gas have been extensively presented in the literature (e.g., Bzowski et al. 2015; Möbius et al. 2015b) and will not be repeated here. In brief, *IBEX* is a spin-stabilized Earth satellite in a very elongated orbit (McComas et al. 2009). To maintain the spin axis within a few degrees from the Sun, it was adjusted once per orbit during the first three years of operation, and is adjusted twice per orbit after the *IBEX* orbit change in 2011 (McComas et al. 2011). The boresight of *IBEX*-Lo points perpendicular to the spin axis. For a given orientation of the spin axis, data are collected from the great circle of the sky visible to the instrument for several days when the spacecraft is sufficiently high above the magnetopause. The registered counts are binned into 6° spin angle bins.

On the ground, the data are filtered against all known perturbations. Due to this filtering, some intervals of the observations are excluded. The intervals used in the analysis are referred to as ISN good times. In this paper, the previously used ISN good times (Leonard et al. 2015; Möbius et al. 2015b) are further restricted by adopting the intervals included in the good times obtained from an in-depth investigation of the sources of background and foreground, presented by Galli et al. (2015, 2016, 2017). This restriction results in shorter accumulation times in several orbits and rejection of all data from orbit 237b. The list of orbits selected for this analysis is presented in Table 1. The background level was also revised

Table 1
List of Orbits Selected for This Analysis

Year	Orbits
2009	14, 15, 16, 17, 18, 19
2010	63, 64, 65, 66, 67
2011	110, ^a 112, ^a 113, ^a 114, ^a 115
2012	153b, ^a 154a, ^a 154b, ^a 156a, ^a 156b, 157a, 157b, 158a
2013	193a, 193b, 194a, 194b, 195a, 195b, 196a, 196b, 197a, 197b
2014	233b, 234a, 234b, 235a, 235b, 236a, 236b, 237a, 238a
2015	274a, 274b, 275a, 275b, 276a, 276b, 277b

Note.

^a Data only from ESA step 2 (ESA steps 1 and 3 not observed due to high-resolution mode).

following results from Galli et al. (2016, 2017). The data subset used now is cleaner, which results in a slightly lower χ^2 obtained in the present fitting than that by Bzowski et al. (2015) as will be evident from the results presented below.

The data used in this analysis include (1) intervals of the good times and (2) the numbers of counts due to ISN atoms, collected during these time intervals in each individual 6° spin angle bin. From this information, we calculate mean counting rates for each bin for a given orbit.

An extensive presentation of the data used for the ISN He analysis is available in Bzowski et al. (2015), and some important aspects of the data processing carried out before these data are used in the analysis were discussed by Möbius et al. (2015a, 2015b) and Swaczyna et al. (2015). The basis for fitting the parameters are orbit-averaged counting rates as a function of spin angle bins. For illustration purposes, however, in Figure 1 we show them as sky maps in the spacecraft inertial frame in ecliptic coordinates centered at the inflow direction ($255.7^\circ, 5^\circ 1'$), averaged over all observation seasons 2009 through 2015. This centering allows for direct comparison of the presented maps with the other maps of the atom fluxes observed by *IBEX* (e.g., McComas et al. 2017). The count rate in each pixel is a weighted average of the 6° spin angle bins that fall into the pixel with the weights resulting from the Poisson uncertainty. The reduced instrument sensitivity after 2012 is compensated by factors of 2.2, 2.2, and 2.1 for ESA steps 1, 2, and 3, respectively (see Section 3). The maps are presented separately for the three ESA steps. In these maps, we mark the data range used for the present analysis of ISN He, as well as the range used by Kubiak et al. (2016) in their analysis of the Warm Breeze component. The data used in the Warm Breeze fitting were taken between mid-November of the preceding year and mid-January. The data collected afterwards until the end of February are used in the ISN He analysis. The range of the data presented in the maps cover only the ISN seasons and thus spread only over part of the sky. The black region (longitudes outside the range $\sim(145^\circ, 270^\circ)$ occupied by the ISN emission) are observed by *IBEX-Lo* but the data are not presented here because they only have contributions from inner heliosheath ENAs (Energetic Neutral Atom; Fuselier et al. 2012; Schwadron et al. 2014; Galli et al. 2017).

It is important to realize that solar gravity significantly modifies the trajectories of ISN He atoms. In addition, *IBEX* is moving relative to the Sun and the ISN atoms it observes, thereby requiring a frame transformation for the observed velocities. Therefore, the direction of entry of an atom into the instrument collimator significantly differs from the direction of

entry of the atom into the heliosphere. The mean magnitudes of this angular shift are different for the primary ISN He and the Warm Breeze populations, as illustrated in Figure 1. The angle between the peak flux observed by *IBEX* and the inflow direction at the heliopause is $\sim 30^\circ$ for the primary ISN population and $\sim 54^\circ$ for the Warm Breeze population, which is indicated by the correspondingly colored circles and dots in this figure.

The observed ISN He flux distribution in the sky, as represented by the measured counting rate, is a result of a convolution of the angular distribution of the atoms entering the instrument with the collimator transmission function of *IBEX-Lo* (Equation (32) in Sokół et al. 2015b). Differences between maps obtained from individual observation campaigns result from several factors. (1) The ionization rate, and thus ionization losses, vary with solar activity. (2) The observation geometry varies from season to season (the spin axis orientations do not precisely repeat from one year to another, and the distribution of good times for equivalent orbits varies between observation seasons). (3) The *IBEX* motion around the Earth varies slightly for equivalent orbits, and additionally there was an important change of the *IBEX* orbit in 2011 (McComas et al. 2011). All those elements are taken into account in the modeling of the *IBEX* signal.

Inspection of the three panels of Figure 1 shows that, generally, the observed signal is comparable in shape and magnitude in all three ESA steps, but there is a clear, systematic reduction of the area in the sky occupied by the signal from lower to higher ESA steps. This effect is especially pronounced in the region occupied solely by the Warm Breeze, i.e., to the right of the ISN signal peak. To the left of the region used for analysis of ISN He, there is a region where a relatively strong signal is visible in the lowest ESA step, while the intensity decreases in ESA steps 2 and 3. This region is believed to be occupied mostly by ISN H and, as shown in the figure, has not been used in the analysis (Saul et al. 2013). How much ISN H contributes to the signal has not been fully determined to date and is a subject of ongoing research.

The total distribution of ISN H is strongly modified in the outer heliosheath due to charge exchange between ISN H and the local disturbed plasma, so the temperatures and bulk velocity vectors of ISN H are different from those of ISN He already at the entrance to the heliosphere (for in-depth discussion of these effects, see Katushkina et al. 2015). Inside the heliosphere, hydrogen atoms follow complex trajectories due to the action of solar radiation pressure. Radiation pressure is a function of the radial velocity of the atom due to the Doppler effect and of the magnitude of the solar Ly α flux, which varies over the solar cycle (Tamopolski & Bzowski 2009). Hydrogen atoms are subject to much stronger ionization losses than helium, and the ionization rate varies differently with time and heliolatitude than that for He. All these effects result in a different location of the peaks of the counting rates due to ISN H and ISN He in *IBEX-Lo* sky maps, and in different ratios of the H/He fluxes in different yearly observation campaigns. In particular, during some seasons (especially those close to the maximum of solar activity) ISN H may be practically unobservable for *IBEX-Lo* (Saul et al. 2013; Galli et al. 2017).

Before fitting the parameters of the ISN He gas, it was necessary to subtract the Warm Breeze contribution from the signal. Even though this contribution is small within the data

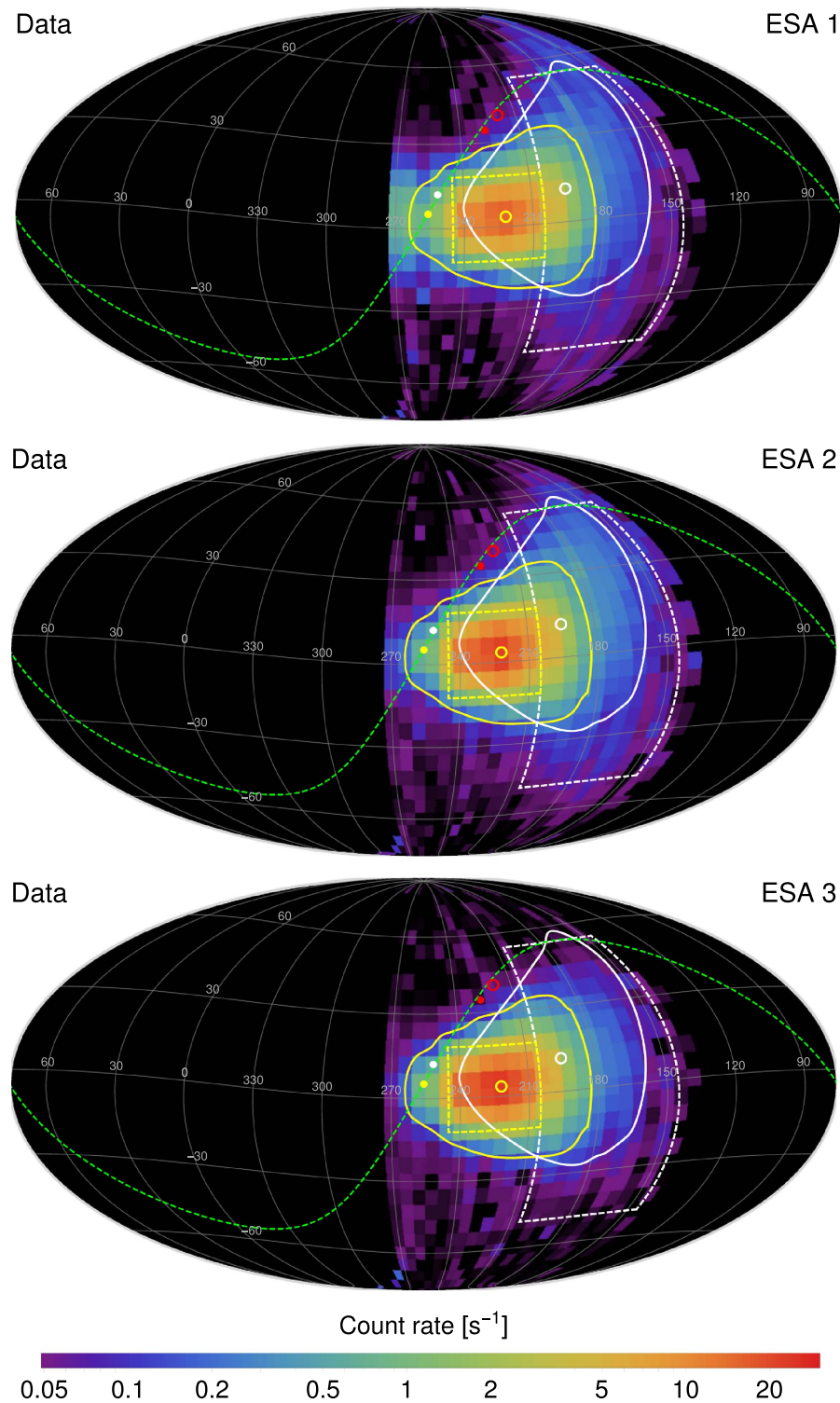


Figure 1. Sky maps of the count rate in s^{-1} due to ISN atoms observed by *IBEX-Lo* in ESA steps 1 (top), 2 (middle), and 3 (bottom) in the J2000 ecliptic coordinates centered at the nose direction in the *IBEX*-inertial frame, averaged over the ISN campaigns 2009 through 2015 (the sensitivity decrease after 2012 is compensated for; see the text). The yellow and white dashed rectangles on the celestial sphere mark the data range used in the fitting of the primary ISN He in Bzowski et al. (2015) as well as in the present paper, and in the fitting of the Warm Breeze in Kubiak et al. (2016), respectively. The solid contours mark where the contribution of the primary ISN He (yellow) and the Warm Breeze (white) signal in ESA step 2 is $\approx 0.1 \text{ s}^{-1}$; note that these contours are identical in all three panels to facilitate viewing the differences between the flux observed in the three ESA steps. The maxima of the two components (marked with open yellow and white circles) fall inside their respective contours. The dashed green line is the deflection plane of the secondary population as obtained by Kubiak et al. (2016). The unperturbed inflow directions of these two populations are shown as the solid yellow and white dots on this plane. The red circle represents the center of the *IBEX* ribbon (Funsten et al. 2013) and the red dot is the direction of the unperturbed interstellar magnetic field, derived by Zirnstein et al. (2016) from global models of the heliosphere using the geometry of the ribbon as constraints.

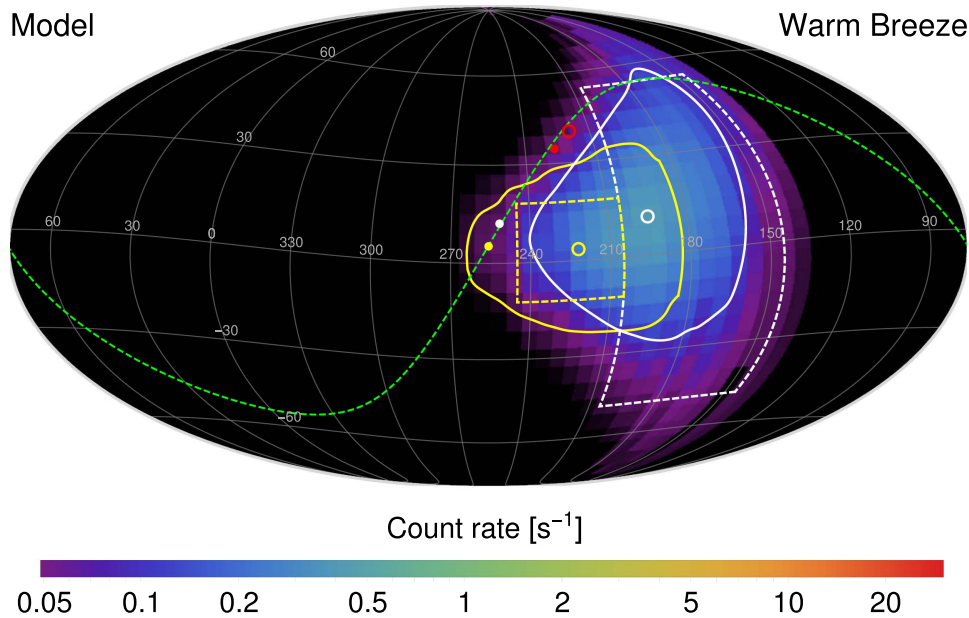


Figure 2. Simulated count rate due to the Warm Breeze in ESA step 2, calculated using the Warm Breeze inflow parameters from Kubiak et al. (2016), averaged over all ISN observation campaigns used in the present analysis. The lines, points, and contours are identical to those in Figure 1.

region adopted for analysis its presence would significantly affect the accuracy of the derivation of the ISN He inflow parameters (Swaczyna et al. 2015). To that end, we calculated the expected count rate due to the Warm Breeze using the model parameters obtained by Kubiak et al. (2016), separately for each of the ISN observation campaigns. The time-averaged contribution from the Warm Breeze is shown in Figure 2. To account for the differences between the sensitivities in the ESA steps used, the model count rate calculated for ESA step 2 was rescaled to steps 1 and 3 by factors 0.8 ± 0.2 and 1.0 ± 0.3 , respectively. These factors were assessed based on the comparison of the observed Warm Breeze signals in the respective ESA steps. In this analysis we do not account for the energy-dependent response in the Warm Breeze. This variation may be even more important for this population than for ISN He but we leave it for future studies. Fortunately, the speed range of Warm Breeze atoms in the range used in the spin angle range used in this analysis is small (see Figure 3). Because of the fit uncertainty of the Warm Breeze parameters, subtracting this model from the data modifies the uncertainty system of the data. This additional uncertainty was accounted for in the parameter fitting.

2.2. Accounting for the Energy Characteristics of IBEX-Lo

Fitting the data and inferring the flow parameters of ISN He is performed using forward modeling. In this approach, a model of the observed gas is calculated for an assumed set of ISN and instrument parameters, and the results are compared with the data. This modeling is done using the Warsaw Test Particle Model (WTPM, Sokół et al. 2015b), which assumes a spatially homogeneous flow of Maxwell-Boltzmann gas everywhere beyond the heliopause. Inside the heliopause, the atoms follow heliocentric hyperbolic orbits and are subject to time-varying ionization losses due to solar photoionization and solar wind particles (electrons and protons), with the rates taken from Bzowski et al. (2013), Sokół & Bzowski (2014), and Sokół et al. (2016).

The model used in the fitting adopts certain assumptions on the efficiency of converting the flux of atoms entering the instrument into a counting rate returned by this instrument. In the studies of the primary population of ISN He by Bzowski et al. (2015) it was assumed that within ESA step 2 there is no energy dependence of the instrument sensitivity, but that the overall sensitivity could change from one observation season to another. The sensitivity level for each season was a free normalization parameter of the fit, further on referred to as the Y-norm (for “yearly norm”). Based on the fit results it was demonstrated that the yearly sensitivities agreed with each other within the time intervals of identical PAC voltage in the instrument. Statistically, this result could be interpreted as no change in the sensitivity between individual seasons other than an approximately twofold reduction in the sensitivity due to the reduction of the PAC voltage in 2012.

Based on the conclusion of the constancy of Y-norms within the intervals of identical PAC values, two normalization parameters for the two observation intervals with the two PAC values were adopted in the analysis of the Warm Breeze by Kubiak et al. (2016). Norms of this kind will be further on referred to as P-norms (for “PAC voltage norms”).

In the present analysis, we used data from three energy steps instead of one and therefore we had to modify the approach to the energy sensitivity in the modeling. In reality, the energy response function for a given species depends on several factors: (1) sputtering efficiency, (2) energy distribution of the sputter products, (3) transmission efficiency for the sputter products through the ESA, and (4) detection efficiency in the time-of-flight section of the detector. Some of these aspects are not fully understood for the actual in-flight environment. Factors (1) and (2) are independent of setting the instrument to a given ESA step because observations in different ESA steps are carried out in close succession, and the flux and speed distribution of the incoming atoms do not change over short times. As a result, the signal observed in ESA steps 1, 2, and 3 is created by the same flux of He atoms. Consequently, the mean counting

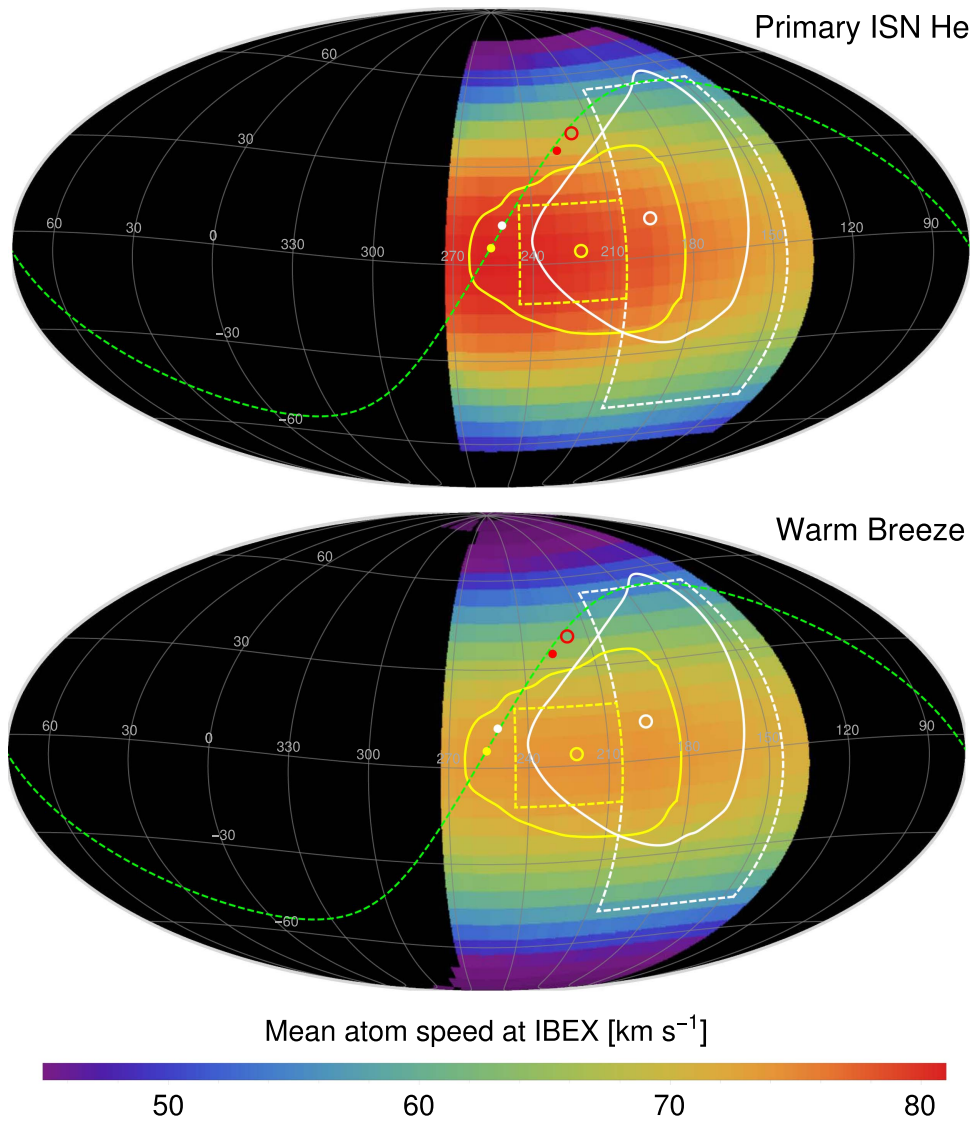


Figure 3. Sky maps of the mean speed of He atoms from the primary population (upper panel) and the Warm Breeze (lower panel) relative to *IBEX*. Different ecliptic latitudes correspond to different spin angle bins. The absolute speed scale in km s^{-1} is shown with the color bar.

rates obtained in a given spin angle bin in a given *IBEX* orbit in these ESA steps differ solely due to the different sensitivities of the instrument in these ESA steps (excluding the inevitable statistical scatter).

From modeling, We found the range of the mean impact velocities of the observed ISN He atoms in each spin angle bin. The resulting mean impact velocities are presented in Figure 3. In this figure, we show that each spin angle bin has its own quite narrow range of atom speeds. Within the range of the data used for the ISN and Warm Breeze analyses, the mean impact velocities for the ISN He atoms range from 78 to 81 km s^{-1} , and for the Warm Breeze from 60 to 75 km s^{-1} . Therefore this unknown sensitivity function can be approximated by a linear function of speed over a relatively narrow speed range. The parameters of these linear relations, different for each ESA step and for each PAC voltage, will be free parameters in the fits. Within the data range used, the atom energies are $\sim 127\text{--}137 \text{ eV}$, well above the threshold energy, discussed in the Introduction, and thus the issue of the magnitude of this threshold speed does not affect our analysis.

2.3. Scaling Functions

The model we adopted in the simulations is an extension of the version of the WTPM discussed in detail by Sokół et al. (2015b). The calculation proceeds identically as discussed in that paper up to Equation (44), which defines the particle scalar flux $F_i(\boldsymbol{\pi})$ in a given *IBEX* orbit, averaged over the width of the spin angle bin and good time intervals. Index i corresponds to the index of a given data point, and $\boldsymbol{\pi}$ is the ISN He parameter set chosen to calculate the given model (the inflow velocity vector and the temperature). Alongside this calculation, we now calculate the mean impact speed for He atoms for each bin $v_i(\boldsymbol{\pi})$. The difference in the calculation of the count rate that we introduce in comparison with Sokół et al. (2015b) and Swaczyna et al. (2015) is the more general scaling function $S(i, v, \boldsymbol{q})$, used here instead of the constant scaling factor adopted previously. The scaling function (expressed in $\text{cm}^2 \text{ sr}$) is a function of the mean atom impact velocity and converts the physical flux F_i (in $\text{cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$) to the simulated count rate g_i (in s^{-1}):

$$g_i(\boldsymbol{\pi}, \boldsymbol{q}) = S(i, v_i(\boldsymbol{\pi}), \boldsymbol{q})F_i(\boldsymbol{\pi}), \quad (1)$$

where v_i is the characteristic speed for a given orbit and spin angle bin, which is now calculated alongside F_i , and \mathbf{q} is a vector of parameters that define the scaling function. Importantly, F_i is a physical characteristic of the atoms observed in this spin angle bin and orbit, which of course does not depend on the ESA step of the instrument. In addition to F_i , the model now returns the average speed and the average squared speed for each bin.

The count rate calculation in Sokół et al. (2015b) was carried out using the scaling function that is constant for each ISN season:

$$S_Y(i, v, \{a_{2009}, a_{2010}, \dots, a_{2014}\}) = a_{y(i)}, \quad (2)$$

where $y(i)$ give the years that represent each of the ISN seasons. This form represents Y-norms and means that the scaling function does not depend on the atom impact speed. Inclusion of more ESA steps would require three separate constants for each ESA step and each ISN season. However, we verified that this is not necessary because the norms for individual seasons are constant within two periods: before and after the decrease of the PAC voltage. Consequently, we define the P-norms using the scaling function in the form:

$$S_P(i, v, \{a_{p,ESAe}\}_{p=H,L;e=1,2,3}) = a_{p(i),ESAe(i)}, \quad (3)$$

where $p(i)$ gives H or L for high/low PAC voltage, depending if a given data point i was obtained before or after the decrease of the PAC voltage, and $e(i)$ gives the ESA step. Using the approach with P-norms implies that a different parameter a is allowed for each energy step and the PAC voltage. As a result, for three energy steps the vector \mathbf{q} of the parameters of function S_P has six elements.

If one allows for a true dependence of the sensitivity of the instrument within a given ESA step on the atom speed, then the simplest possible form of the scaling function is the linear function S_{PV} , defined as:

$$S_{PV}(i, v, \{a_{p,ESAe}, b_{p,ESAe}\}_{p=H,L;e=1,2,3}) = a_{p(i),ESAe(i)}(1 + b_{p(i),ESAe(i)}(v - v_0)), \quad (4)$$

where v_0 is a certain reference speed, adopted here at 78 km s^{-1} . This arbitrary choice does not affect the generality of the scaling function because it only requires a simple transformation for the parameters $a_{p,ESAe}$ for a different choice of v_0 . The function S_{PV} may be regarded as an expansion of an unknown true sensitivity function into a Taylor series, cut off at the linear term. The parameters $b_{p,ESAe}$ are the relative change of the absolute sensitivity per km s^{-1} . In general one may introduce a function S_{YV} defined in analogy to S_{PV} but with the parameters depending on the season, and not the H/L PAC voltage.

2.4. Parameter Fitting

Parameter fitting is carried out using the method suggested by Swaczyna et al. (2015) and successfully applied by Bzowski et al. (2015) to obtain the inflow parameters of ISN He and by Kubiak et al. (2016) to obtain the inflow parameters of the Warm Breeze. Swaczyna et al. (2015) discuss at length all data correlations that must be taken into account and we will not repeat all those details. We only briefly recall that the model must be calculated on a regular parameter grid and that one

must compute the uncertainty covariance matrix \mathbf{V} . With those on hand, we calculated the measure of goodness of fit given by χ^2 defined as:

$$\chi^2(\boldsymbol{\pi}, \mathbf{q}) = \sum_{i,j} (c_i - g_i(\boldsymbol{\pi}, \mathbf{q}))(\mathbf{V}^{-1})_{ij}(c_j - g_j(\boldsymbol{\pi}, \mathbf{q})) \quad (5)$$

where i, j go over all data points used in the analysis, for all three ESA steps, c_i is the corrected count rate with the background and Warm Breeze contribution subtracted (Swaczyna et al. 2015), and g_i is calculated from Equation (1) with the S functions given by Equations (2)–(4). The fit parameters are $\boldsymbol{\pi}$ and \mathbf{q} . The parameters $\boldsymbol{\pi}$ correspond to the sought ISN parameters (the magnitude and ecliptic coordinates of the inflow velocity, and the temperature), and the parameters \mathbf{q} are the unknown parameters of the sensitivity function S . The calculation was performed on sets of the ISN He parameters taken from a grid spaced at $\Delta\lambda = 1^\circ$, $\Delta\beta = 0.1$, $\Delta v = 0.2 \text{ km s}^{-1}$, $\Delta v_t = 0.1 \text{ km s}^{-1}$, with 4607 nodes around the ISN He parameter values obtained by Bzowski et al. (2015).

The optimum parameter set was found in a two-step process. The first step was finding the parameter set $(\boldsymbol{\pi}, \mathbf{q})$ with the minimum χ^2 value calculated from Equation (5), evaluated for those 4607 parameter sets. Subsequently, a quadratic form in the parameter space was fitted to the χ^2 surface near the intermediate minimum, and the parameter values for the absolute minimum of this surface were calculated analytically. This resulted in a better parameter resolution than that obtained directly from the parameter grid spacing. The fitting procedure is repeated for each considered data set and scaling function form. Consequently, both the inflow parameters and the parameters \mathbf{q} vary accordingly.

Statistically, the expected value of χ^2_{\min} is given by $\chi^2_{\min} \simeq N_{\text{dof}} \pm \sqrt{2N_{\text{dof}}}$, where N_{dof} is the number of degrees of freedom. For N_{dof} we adopted the number of data points minus the number of fit parameters, similarly to Swaczyna et al. (2015), Bzowski et al. (2015), and Kubiak et al. (2016), even though this is strictly exact only for linear models, and our model is nonlinear. Finally, the uncertainties of the sought parameters are obtained from the curvature of the function $\chi^2(\boldsymbol{\pi}, \mathbf{q})$ scaled using the scaling factor $S = \sqrt{\chi^2/N_{\text{dof}}}$ as presented in Swaczyna et al. (2015) and Bzowski et al. (2015). The parameters $(\boldsymbol{\pi}, \mathbf{q})$ are correlated with each other, as discussed in Swaczyna et al. (2015).

3. Results

We began the parameter evaluation by repeating the analysis of Bzowski et al. (2015). We used a slightly more restrictive good times list and slightly modified levels of the ubiquitous background (Galli et al. 2016, 2017). The contribution from the Warm Breeze was estimated from WTPM using the Warm Breeze parameters from Kubiak et al. (2016), i.e., different from the Warm Breeze estimate used by Bzowski et al. (2015). This contribution was subtracted from the data and the related uncertainty added to the uncertainty system. We performed fits assuming Y-norms and using data from observation seasons 2009–2014 and ESA step 2 only. The ISN He fit parameters are presented in row 1 in Table 2, and the corresponding Y-norms in Table 3. In rows 0 of these tables the parameters found by Bzowski et al. (2015) are presented for comparison. We conclude that the new data selection, as well as the new Warm

Table 2
ISN He Flow Parameters

	Norm	ESA	Seasons	λ ($^\circ$)	β ($^\circ$)	v (km s $^{-1}$)	T (K)	M	N_{dof}	$\frac{\chi^2}{N_{\text{dof}}}$
0 ^a	Y-norm	2	2009–2014	255.74 \pm 0.45	5.16 \pm 0.10	25.76 \pm 0.37	7440 \pm 260	5.079 \pm 0.028	254	1.84
1	Y-norm	2	2009–2014	255.70 \pm 0.47	5.14 \pm 0.11	25.71 \pm 0.43	7696 \pm 294	4.980 \pm 0.020	248	1.60
2	Y-norm	2	2009–2015	255.68 \pm 0.46	5.09 \pm 0.11	25.65 \pm 0.39	7677 \pm 285	4.976 \pm 0.019	289	1.55
3	P-norm	2	2009–2015	255.46 \pm 0.45	5.11 \pm 0.11	25.85 \pm 0.40	7805 \pm 278	4.972 \pm 0.019	294	1.61
4	PV-norm	2	2009–2015	255.48 \pm 0.46	4.91 \pm 0.09	26.23 \pm 0.42	7726 \pm 269	5.072 \pm 0.025	292	1.42
5	P-norm	2–3	2009–2015	254.76 \pm 0.33	5.20 \pm 0.08	26.49 \pm 0.30	8111 \pm 211	4.999 \pm 0.015	544	1.65
6	PV-norm	2–3	2009–2015	254.90 \pm 0.36	5.04 \pm 0.08	26.66 \pm 0.33	7938 \pm 208	5.085 \pm 0.019	540	1.45
7	P-norm	1–3	2009–2015	254.63 \pm 0.35	5.23 \pm 0.07	25.98 \pm 0.29	7853 \pm 208	4.982 \pm 0.016	794	1.97
8	PV-norm	1–3	2009–2015	253.82 \pm 0.37	5.06 \pm 0.07	26.98 \pm 0.33	8155 \pm 215	5.077 \pm 0.020	788	1.74
9	P-norm	1–3	2009–2015 ^b	255.62 \pm 0.36	5.16 \pm 0.08	25.82 \pm 0.33	7673 \pm 225	5.010 \pm 0.015	686	1.64
10	PV-norm	1–3	2009–2015 ^b	255.41 \pm 0.40	5.03 \pm 0.07	26.21 \pm 0.37	7691 \pm 230	5.080 \pm 0.019	680	1.49

Notes.^a Results from Bzowski et al. (2015).^b Data range restricted to longitudes within 115 $^\circ$ –155 $^\circ$ (see the text).**Table 3**
Y-norms Fitted along with the ISN Inflow Parameters ($a_{y,\text{ESA}2}$, 10 $^{-6}$ cm 2 sr)

	2009	2010	2011	2012	2013	2014	2015
0 ^a	16.76	16.66	15.11	14.50	7.709	7.977	...
1	17.18 \pm 0.19	17.08 \pm 0.20	17.47 \pm 0.28	17.27 \pm 0.22	8.03 \pm 0.10	8.25 \pm 0.10	...
2	17.16 \pm 0.19	17.07 \pm 0.20	17.45 \pm 0.28	17.25 \pm 0.22	8.02 \pm 0.09	8.25 \pm 0.10	8.28 \pm 0.09

Note. The row numbers correspond to the cases presented in Table 2.^a Results from Bzowski et al. (2015).

Breeze subtraction, does not substantially affect the ISN He inflow parameters. The differences are small, well within the uncertainties. The largest discrepancy is in the temperature (\sim 250 K), and this is likely caused by the new Warm Breeze model subtracted here. The temperature of the Warm Breeze was found significantly lower by Kubiak et al. (2016) than obtained from the earlier analysis by Kubiak et al. (2014). Bzowski et al. (2015) adopted the higher temperature of the Warm Breeze from Kubiak et al. (2014). Consequently, after subtraction of the Warm Breeze in their analysis, the data featured a narrower peak, which the fitting procedure interpreted as due to a lower temperature of ISN He. The reduced χ^2 value that we have obtained now, equal to χ^2/N_{dof} , is significantly lower than that obtained by Bzowski et al. (2015), i.e., the model better fits the data than in Bzowski et al. (2015).

In the next step, we added the data from season 2015 to the analysis (row 2). The resulting parameters remain almost the same. Inspection of Table 3 shows that the obtained parameters of the scaling function are constant for the seasons with the H PAC voltage and separately with the L PAC voltage. The variation of these parameters is reduced by an order of magnitude with respect to that obtained by Bzowski et al. (2015). Moreover, the new values agree in these two periods, and the usage of the P-norms rather than Y-norms is justified. Results for the P-norms applied to ESA step 2 solely are presented in row 3 of Tables 2 and 4. The obtained values of the norms agree with the values formerly found for each season. The ISN flow parameters are slightly different, but well within the uncertainty ranges.

We found that including data from the 2016 campaign significantly increases the value of the reduced χ^2 . During this season, only data from three orbits before the absolute

maximum of ISN He count rate could be obtained due to a temporary issue with the *IBEX* star tracker system, which makes the data coverage for this campaign significantly poorer. Therefore, we decided to not include the data from this season in this analysis.

Next, we checked how the fit parameter values change when a variation of the instrument sensitivity with atom speed is allowed for, while continuing to use data solely from ESA step 2. We used the function S_{PV} from Equation (4) and list the results in row 4 of Table 2. This modification resulted in a significant reduction of the reduced χ^2 value. The inflow latitude and speed changed appreciably but within the uncertainty range.

Subsequently, we tested whether including additional ESA steps while assuming no dependence of the instrument sensitivity to atom speed within individual ESA step returns similar results to those obtained by Bzowski et al. (2015), using again the sensitivity function S_{P} . Finally, we allowed for the sensitivity within an ESA step to vary with atom speed while continuing to use data from three ESA steps. To that end, we took PV-norms and adopted the sensitivity function S_{PV} from Equation (4). The resulting ISN He parameters are presented in rows 5–8 of Table 2, and the norms in Table 4. As a result of these fittings we found that the ISN flow parameters are substantially different if all three ESA steps are included (rows 7 and 8) from those obtained solely for ESA step 2, especially for the inflow longitude. This effect is not present if ESA step 1 is excluded (rows 5 and 6). The simultaneous increase in the reduced χ^2 value for the three ESA steps suggests that the model fit is significantly poorer. This is likely due to a contribution from ISN H, unaccounted for in the model. We consider this hypothesis is likely because eliminating ESA step 1 from the data and fitting to data from ESA steps 2 and 3

Table 4
P-norms and PV-norms Fitted along with the ISN Inflow Parameters

	H, ESA1	L, ESA1	H, ESA 2	L, ESA 2	H, ESA 3	L, ESA 3
Parameters $a_{p,ESAe}$ (10^{-6} cm ² sr)						
3	17.27 ± 0.17	8.12 ± 0.09
4	19.26 ± 0.38	8.63 ± 0.14
5	17.44 ± 0.14	8.25 ± 0.06	20.78 ± 0.19	10.22 ± 0.08
6	19.63 ± 0.36	8.80 ± 0.16	22.06 ± 0.50	10.35 ± 0.20
7	14.14 ± 0.14	6.29 ± 0.05	17.20 ± 0.14	8.18 ± 0.06	20.56 ± 0.20	10.12 ± 0.08
8	14.66 ± 0.41	7.09 ± 0.16	20.60 ± 0.41	9.31 ± 0.18	22.85 ± 0.58	10.96 ± 0.23
9	14.08 ± 0.14	6.27 ± 0.06	17.22 ± 0.14	8.13 ± 0.07	20.50 ± 0.19	10.09 ± 0.08
10	14.49 ± 0.34	6.57 ± 0.13	19.20 ± 0.35	8.63 ± 0.15	21.32 ± 0.49	10.20 ± 0.19
Parameters $b_{p,ESAe}$ (10^{-2} km ⁻¹ s)						
4	-5.8 ± 0.8	-3.7 ± 0.8
6	-5.6 ± 0.7	-4.0 ± 0.8	-3.3 ± 0.9	-1.6 ± 0.9
8	-1.9 ± 1.0	-5.5 ± 0.9	-4.2 ± 0.9	-6.0 ± 0.8	-4.2 ± 0.9	-3.9 ± 0.9
10	-1.8 ± 1.0	-3.1 ± 1.0	-5.6 ± 0.8	-3.8 ± 0.9	-2.3 ± 1.0	-1.2 ± 1.0

Note. The row numbers correspond to the cases presented in Table 2.

results in parameters similar to those obtained previously (i.e., no important change compared to Bzowski et al. 2015), and our analysis shows that the presence of ISN H is the most prominent in ESA step 1, consistently with the findings by Saul et al. (2012) and Schwadron et al. (2013).

Most of the ISN H signal can be avoided if one adopts a tighter restriction on the spin axis orientation, which corresponds to the ecliptic longitude of *IBEX*. Here, we limited them to longitudes in the range (115° , 155°) compared to (115° , 160°) used so far. Effectively, this removed data from those orbits where the expected contribution from ISN H was the largest within the data set. As a result, we noticed that this restriction improved the fit quality as it reduced the reduced χ^2 value by ~ 0.3 both for the case with P-norms and for that with PV-norms. This restriction effectively resulted in leaving out six late orbits, i.e., those well after the peak of ISN He is observed (orbits 19, 67, 115, 157b, 158a, 238a). The contribution from ISN H to the signal observed during these orbits is expected to be the largest during a given ISN season. Therefore, as the final result for the ISN He parameters we adopt those listed in rows 9 and 10 in Table 2, with the parameters of the scaling functions in Table 4.

The obtained parameters of the scaling functions suggest that the efficiency of *IBEX-Lo* is increasing within the three ESA steps. Moreover, the high-to-low PAC voltage ratios of the parameters $a_{H,ESAe}/a_{L,ESAe}$ are similar for all three ESA steps (2.21, 2.22, and 2.09 for ESA steps 1, 2, and 3, respectively, from the case presented in row 10 in Table 4).

4. Discussion and Conclusions

The differences between the present analysis of the bulk velocity and temperature of ISN He and that carried out by Bzowski et al. (2015) are the following: (1) subtracting from the data of a better model of the Warm Breeze (from Kubiak et al. 2016 instead of that from Kubiak et al. 2014); (2) adopting more stringent criteria for good times and using slightly different background level for the observation seasons after 2012 (Galli et al. 2016); (3) using P-norms instead of Y-norms (i.e., adopting instrument sensitivity parameters characteristic for a given PAC voltage rather than for a given observation season); (4) using data from one more ISN

observation season (2015) while cutting off some late orbits during each season; (5) using data from ESA steps 1 through 3 instead of only ESA step 2; (6) allowing for the instrument sensitivity to depend on the impact speed of He atoms within individual ESA steps (function S_{PV} from Equation (4) instead of a constant value).

Modifications (1) and (2) result in a small change in the Mach number of the flow, reducing it from 5.079 ± 0.028 obtained by Bzowski et al. (2015) to 4.980 ± 0.020 . This reduction is mostly due to an increase in the fitted temperature by 230 K. While not important for the global image of the heliosphere, this is a statistically significant difference in the Mach number of more than 3.5σ . Modifications of the direction and speed of inflow and of the temperature of the ISN He are within $1-2\sigma$. In this respect, switching from Y-norms to P-norms and using data from one more observation season (modifications 3 and 4), as well as using data from three ESA steps instead of one (5), do not result in any further statistically significant modifications of the ISN He parameters. Adding more data and applying more stringent filtering results in a certain reduction of χ^2 per degree of freedom, but still these χ^2 significantly exceed the expected value. The small increase in the temperature is understandable since the temperature of the Warm Breeze obtained by Kubiak et al. (2016) was lower than that obtained by Kubiak et al. (2014). As a result, the change in the modeled Warm Breeze signal subtracted from the data was the largest for the bins at spin angles farthest from the peaks, which resulted in a larger signal in these bins left in the ISN He signal used in parameter fitting.

Allowing for the instrument sensitivity within an ESA step to vary with the impact speed (modification 6) restores the Mach number to the value originally obtained by Bzowski et al. (2015). Also, the other parameters are little affected in comparison with those obtained by Bzowski et al. (2015), even though the latitude of the flow is slightly lower, and the temperature and the speed are slightly larger. From this analysis we conclude that the current best estimate for the bulk velocity vector and temperature of ISN He are those listed in the last two rows of Table 2. These two parameter sets are in good agreement. The difference between them is presented in Table 5. This table shows also a comparison with the results

Table 5
ISN He Flow Parameters from Different Studies of *IBEX* Data

Reference	λ ($^{\circ}$)	β ($^{\circ}$)	v (km s $^{-1}$)	T (K)
This analysis P-norm (9)	255.62 ± 0.36	5.16 ± 0.08	25.82 ± 0.33	7673 ± 225
This analysis PV-norm (10)	255.41 ± 0.40	5.03 ± 0.07	26.21 ± 0.37	7691 ± 230
Difference (9)–(10)	0.21	0.13	−0.39	−18
Bzowski et al. (2015)	255.74 ± 0.45	5.16 ± 0.10	25.76 ± 0.37	7440 ± 260
Schwadron et al. (2015)	255.6 ± 1.4	5.12 ± 0.27	25.4 ± 1.1	8000 ± 1300
McComas et al. (2015b) ^a	255.7	5.1	25.4	7500

Note.

^a “Working values” based on Bzowski et al. (2015) and Schwadron et al. (2015).

of the previous studies of the *IBEX* data. The parameters obtained in this analysis are in agreement with the “working values” provided by McComas et al. (2015b) based on the analyses by Bzowski et al. (2015) and Schwadron et al. (2015).

The combination of the model of ISN He obtained now and of the Warm Breeze model from Kubiak et al. (2016) is still not perfect. The most significant issue is likely the presence of ISN H in the data. This hypothesis is based on analysis of residuals, i.e., difference between the count rates in the data and the model. Maps of these residuals, averaged over all observation seasons, are presented in Figure 4 separately for the three ESA steps. The yellow, orange, and red colors mark positive differences, i.e., an excess of the data over the model. As clearly seen, this excess is mostly located in the region where ISN H is expected based on modeling of the ISN H signal (see, e.g., Kubiak et al. 2013). This is also in qualitative agreement with data analyses in Saul et al. (2012, 2013) and Schwadron et al. (2013). Both the magnitude and spatial range of the excess are reduced when going from ESA steps 1–3 and they evolve with time, as is clearly implied from a comparison of the map of the residuals for ESA step 1, ISN observation seasons 2009 and 2010, shown in the upper panel of Figure 5, with those obtained for ISN seasons 2013 through 2015, shown in the lower panel of this figure.

It is important to note that the residual excess is present at all times, which suggests that it is indeed due to the physical reason just discussed and not necessarily to an inadequacy of the adopted model of the ISN He and Warm Breeze. Furthermore, the intensity of the excess qualitatively conforms with the expectations for the evolution of ISN H flux during the solar cycle, with a stronger signal expected during low solar activity and a reduction in the signal strength when the activity increases (Kubiak et al. 2013). In our case, the excess for the ISN seasons 2009 and 2010 is larger than that for the solar maximum seasons 2013–2015. This is understandable because during 2009 and 2010 the Sun had just emerged from a prolonged minimum of activity when both solar radiation pressure acting on interstellar H and the ionization losses were lowest. By contrast, during 2013–2015 ISN seasons the Sun was at the maximum of its activity and both the high ionization rate (see Figure 31 in McComas et al. 2017) and high radiation pressure reduced the flux of ISN H at 1 au considerably. As a result, the flux of ISN H at 1 au during this time is expected to be lower than during the minimum of solar activity.

The magnitude of the residuals likely underestimates the true magnitude of the ISN H contribution to the observed signal. This is because the fitting procedure aims to fit the data to the model as well as possible within the adopted model, which in our case does not include ISN H. The fitting procedure we have

used does not take into account possible systematic patterns in the residuals. In the case of a perfect fit, with all signal components taken into account, the spatial distribution of positive and negative residual values should be random. In our case, evidently connected regions of negative and positive residuals exist. The region of positive residuals extends well beyond the data subset used in the fitting. These are indications that the model, which only includes the primary ISN He population and the Warm Breeze, is missing an extra component, which is likely ISN H. Since, however, the optimization procedure seeks to have the mean value of the residuals close to zero, some counts due to ISN H were inadvertently attributed to ISN He. Therefore it is not recommended to adopt the residuals as a signal from ISN H at face value. The true contributions from ISN H must be assessed separately, based on a model that contains all components ISN H, ISN He, and the Warm Breeze. Moreover, the residuals to the right of the ISN He peak show a cluster of pixels with negative values surrounded by pixels with positive values in ESA steps 1 and 2. The same region shows predominantly negative residuals in ESA step 3. It suggests that the model of the Warm Breeze also needs to be revised.

The presence of neutral H in the signal during all ISN observation seasons suggests that before further analysis of details of ISN He and the Warm Breeze the question of ISN H must be addressed. This topic was discussed by Schwadron et al. (2013) and Katushkina et al. (2015), who pointed out that the ratio of counts from ISN H actually observed in ESA steps 1 and 2 is significantly different from that expected from state of the art models of ISN H in the heliosphere and that the likely culprit is an inadequate knowledge of the solar Ly α radiation pressure, in particular of the spectral profile of this line. Analysis of the largest currently available data set of full-disk solar Ly α line profiles by Kowalska-Leszczynska et al. (2018) (43 profiles measured during various phases of solar activity) showed that, indeed, the central reversal of the profile is deeper than that used by Katushkina et al. (2015), who adopted a model by Tarnopolski & Bzowski (2009), which had been developed based on just nine profiles then available. Analysis of the implications of this finding for the ISN H signal expected from *IBEX*-Lo is underway now.

The question of the sensitivity of *IBEX*-Lo to He atoms in various ESA steps is interesting mostly in the studies of the secondary population of ISN He, i.e., the Warm Breeze. Studying the energy spectrum of the Warm Breeze may potentially bring a better insight into the plasma flow and temperature in the outer heliosheath. Also, as shown by our results, expanding the analysis of ISN He from one to three ESA steps improves the statistics. Additionally, one obtains a

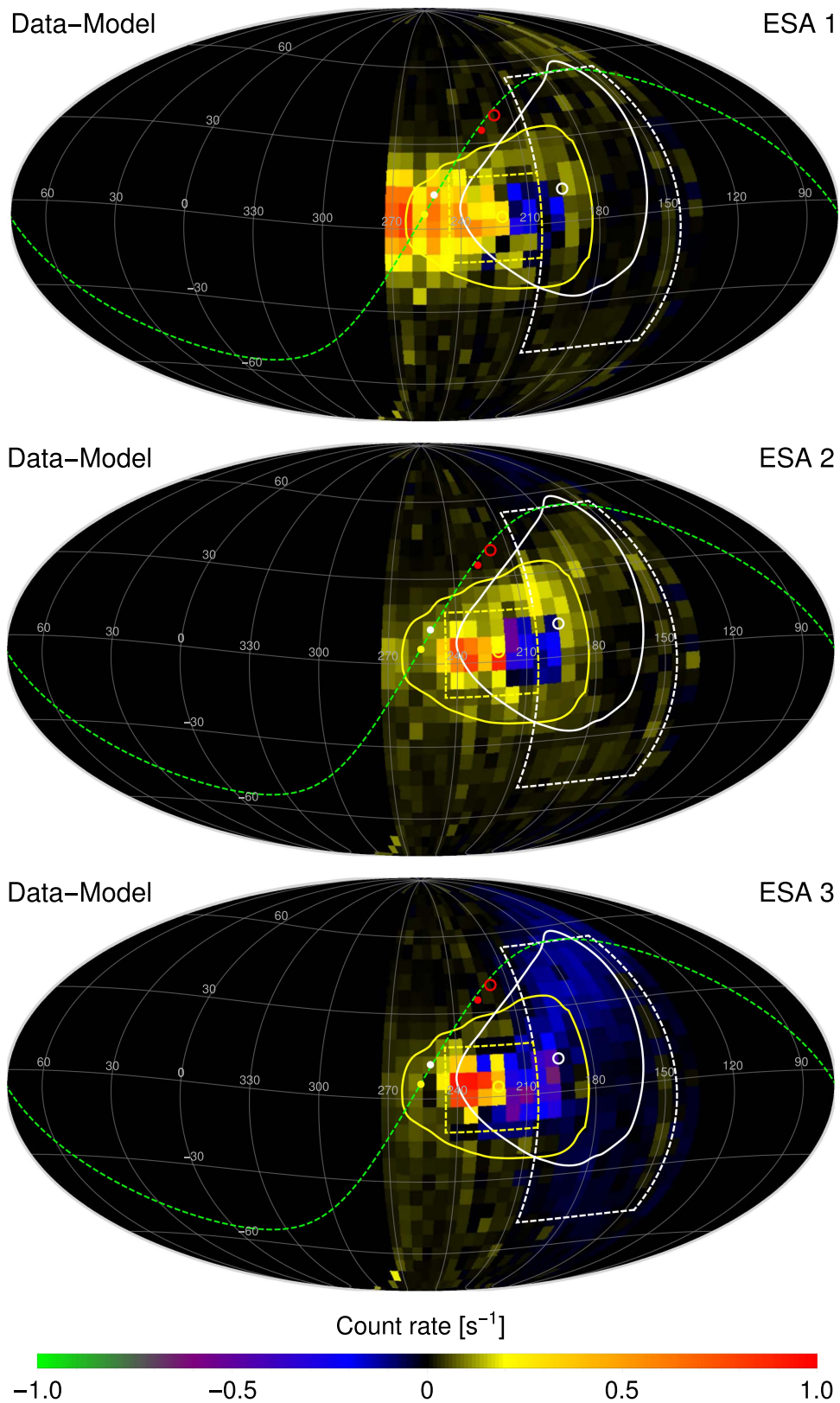


Figure 4. Sky maps of the residuals (data minus model), weighted-averaged over all ISN observation seasons (2009–2015), shown as count rate in s^{-1} in ESA steps 1 (top), 2 (middle), and 3 (bottom) in the J2000 ecliptic coordinates in the *IBEX*-inertial frame. The lines, points, and contours are identical to those in Figure 1.

better opportunity to study ISN H, which is the most abundant in ESA step 1. Therefore understanding the differences in the sensitivity of *IBEX*-Lo to neutral He in the lowest-energy ESA steps is important.

Our analysis supports the conclusion by Bzowski et al. (2015) and Möbius et al. (2015a) that reducing the PAC voltage after the 2012 ISN observation season resulted in a reduction of the overall sensitivity of *IBEX*-Lo to He atoms by

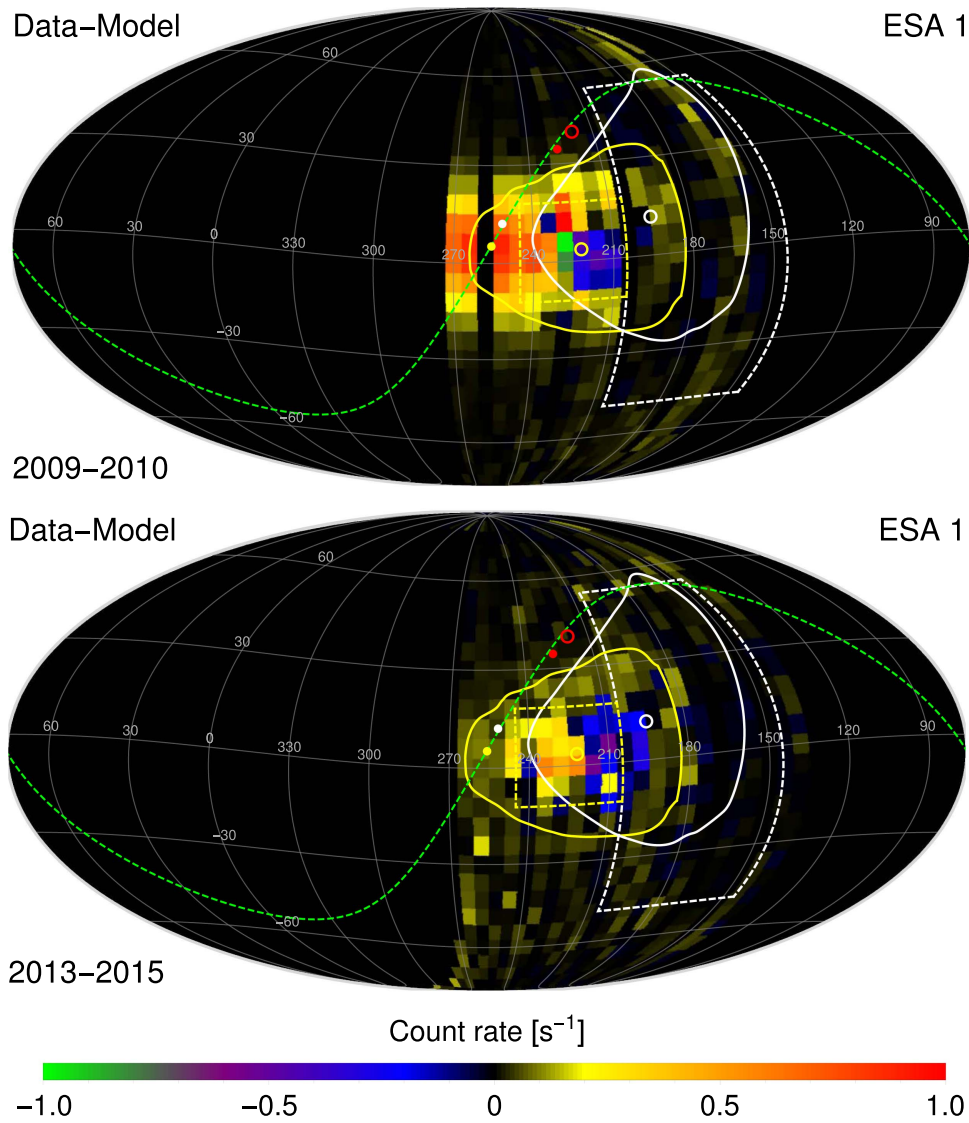


Figure 5. Sky maps of the residuals (data minus model) for ESA step 1, averaged over ISN observation seasons 2009 and 2010 from the minimum of solar activity (upper panel) and over ISN seasons 2013 through 2015 from the maximum of solar activity (lower panel). The lines, points, and contours are identical to those in Figure 1.

factor of ~ 2.2 for ESA steps 1 and 2, and by a factor of ~ 2.1 for ESA step 3. These reductions were taken into account when calculating the averaged maps shown in Figures 1, 4, and 5.

An interesting and surprising finding is that the slope of the speed dependence of sensitivity within the energy steps is negative, i.e., the sensitivity is larger for *slower* speeds of the incoming neutral atoms. This speed dependence came out largest for ESA step 2, consistently for both PAC voltage magnitudes used. The value of $b \sim 5 \times 10^{-2} \text{ km}^{-1} \text{ s}$, typical for ESA step 2 (see Table 4), corresponds to a change of the scaling function by $\sim 25\%$ for the range of the mean velocity of incoming atoms in the selected pixels of $\Delta v \sim 5 \text{ km s}^{-1}$. This behavior is contrary to intuition, especially when one recalls that Kubiak et al. (2014), Galli et al. (2015), and Sokół et al. (2015a) detected a lower energy threshold for the sensitivity in ESA step 2. Whether this is indeed a behavior of the sputtered ion distribution or by how much this effect is produced by the presence of ISN H in the observations, still needs to be evaluated. We believe that either there must be a sensitivity maximum for a certain unknown atom speed within an ESA

step, which is missed due to the adoption of the linear sensitivity functions $S_{pV}(i, v)$ or it is an artifact resulting from the presence of an unaccounted contribution from ISN H in the signal.

It is important to realize that the speed sensitivity within the ESA steps was established only for the largest speed values of neutral He ($75\text{--}81 \text{ km s}^{-1}$), characteristic for the spin angle bins with the largest count rate in the orbits used in the analysis. The other pixels, as well as those occupied by the signal from the Warm Breeze, have lower characteristic speeds (down to $\sim 60 \text{ km s}^{-1}$; see Figure 3). Therefore extrapolating the sensitivity functions resulting from our analysis significantly downward in speed is strongly discouraged. Free from this issue are the parameters obtained assuming that the sensitivity within ESA steps is independent of energies.

Further investigation of the primary and secondary populations of ISN He must be carried out in parallel with analysis of the ISN H component in the data, in an iterative way. As a first step, the contribution from ISN H should be tentatively identified and subtracted from the data. With this, one of the

parameter sets from the fits we have obtained in this paper should be adopted and the model signal due to ISN He should be computed and subtracted from the data. Subsequently, a new parameter set for the Warm Breeze fitted should be fitted to the resulting data set, and the resulting Warm Breeze model used in the next iteration of parameter fitting for ISN He and ISN H.

5. Summary

We have analyzed observations of ISN He carried out by *IBEX* during the ISN He observation seasons 2009 through 2015, for the first time using information from *IBEX*-Lo ESA steps 1 through 3. We have established the differences in the sensitivity to ISN He atoms in these ESA steps and the sensitivity change due to the change in the *IBEX*-Lo PAC voltage setting that was introduced after the 2012 ISN season. We found that the overall sensitivity increases from lower energy steps to the higher ones. Surprisingly, however, we found that the sensitivity within all these ESA steps may be a *decreasing* function of atom speed. We are currently unable to verify if this is a true effect or a result of the presence of counts due to ISN H in the signal.

The excesses of the measured count rates over the modeled signal on the left side of the ISN He peak position in all ESA steps suggests that the data may be partially contaminated by ISN H atoms. We suggest that the signal due to ISN H must be quantitatively interpreted and subtracted before information on the Warm Breeze, available in the measurements in ESA steps 1–3, can be fully utilized. Finally, we found that the bulk velocity vector and temperature of ISN He obtained in the present analysis using data from ESA steps 1–3, listed in the last two rows in Table 2, are within the fit uncertainties of these parameters obtained by Bzowski et al. (2015) from the analysis using only data from ESA step 2. Further analysis of the primary and secondary populations of ISN He and ISN H should be carried out iteratively, using parameters of these populations obtained in a preceding step of the analysis in the subsequent steps.

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